

CHARACTERISTIC FUNCTIONS

We work in 1-dimension (at first)

Defn: $\mu \in \mathcal{P}(\mathbb{R})$. Its characteristic function or Fourier transform, denoted ψ_μ or $\hat{\mu}$ is a function from \mathbb{R} to \mathbb{C}

$$\begin{aligned}\hat{\mu}(t) &= \int_{\mathbb{R}} e^{itx} d\mu(x) \\ &= \int \cos(tx) d\mu(x) + i \int \sin(tx) d\mu(x).\end{aligned}$$

If $x \sim \mu$, we also write ψ_x to mean ψ_μ .

Remarks: Earlier we defined the moment generating function of μ to be $\lambda \rightarrow \mathbb{E}[e^{-\lambda X}]$

which may be represented as $\psi_\mu(-i\lambda)$. For μ supported on $\mathbb{N} = \{0, 1, 2, \dots\}$, the probability generating function $F(t) = \sum_{n=0}^{\infty} \mu_n t^n$ can also be written as $\psi_\mu(-i \log t)$ (at least for $t > 0$).

Thus c.f. is similar to these other 'integral transforms' of μ , but unlike the mgf or pgf, c.f. exists $\forall \mu \in \mathcal{P}(\mathbb{R})$

Basic properties of c.f.: Let $\mu \in \mathcal{P}(\mathbb{R})$. Then

$$\hat{\mu}(0) = 1; |\hat{\mu}(t)| \leq 1 \quad \forall t \in \mathbb{R}; \quad t \rightarrow \hat{\mu}(t) \text{ is uniformly continuous.}$$

Proof: The first two are obvious, since $|e^{itx}| = 1$ ($t \in \mathbb{R}, x \in \mathbb{R}$)

For the last, note that

$$\begin{aligned}|\hat{\mu}(t+h) - \hat{\mu}(t)| &\leq \int |e^{i(t+h)x} - e^{itx}| d\mu(x) \\ &= \int |e^{ihx} - 1| d\mu(x)\end{aligned}$$

As $h \rightarrow 0$, $|e^{ihx} - 1| \rightarrow 0 \quad \forall x$. Further

$|e^{ihx} - 1| \leq 2$ which is integrable w.r.t μ .

By DCT, we get uniform continuity. \square

The importance of characteristic functions stems from the following facts.

(A) Certain operations on measures such as shifting and scaling or convolutions, lead to simple ~~the~~ operations in terms of their characteristic functions.

(B) The characteristic function determines the measure - i.e., $\hat{\mu} = \hat{\nu} \Rightarrow \mu = \nu$.

(This fact alone is sufficient for our purposes, but there are 'explicit' formulas for recovering μ from $\hat{\mu}$, and more usefully, simple formulas to extract useful functionals such as the mean and variance of μ from a knowledge of $\hat{\mu}$).

$$(C) \quad \hat{\mu}_n(t) \rightarrow \hat{\mu}(t) \text{ pointwise } \forall t \in \mathbb{R} \iff \mu_n \xrightarrow{d} \mu.$$

This gives an alternate way to prove weak convergence. It is

sometimes better than working

with CDFs to prove $F_{\mu_n}(t) \rightarrow F_\mu(t)$

\forall continuity points t (for example,

if $\mu_n = \nu_1 * \nu_2 * \dots * \nu_n$, then

F_{μ_n} is very complicated to write in

terms of $F_{\nu_1}, \dots, F_{\nu_n}$, but

$$\hat{\mu}_n(t) = \hat{\nu}_1(t) \hat{\nu}_2(t) \dots \hat{\nu}_n(t) !!)$$

A) Proposition: Let X, Y be real-valued random variables on $(\mathbb{R}, \mathcal{F}, \mathbb{P})$

(a) $\forall \alpha, \beta \in \mathbb{R}, \psi_{\alpha X + \beta Y}(t) = e^{i\beta t} \psi_X(\alpha t)$

(b) If X, Y are independent, then

$$\psi_{X+Y}(t) = \psi_X(t) \psi_Y(t).$$

Proof: (a) $\psi_{\alpha X + \beta Y}(t) = E[e^{it(\alpha X + \beta Y)}]$
 $= e^{i\beta t} E[e^{i\alpha t X}]$
 $= e^{i\beta t} \psi_X(\alpha t).$

(b) $\psi_{X+Y}(t) = E[e^{it(X+Y)}]$
 $= E[e^{itX} \cdot e^{itY}]$
 $= E[e^{itX}] E[e^{itY}]$ (independence)
 $= \psi_X(t) \psi_Y(t). \quad \blacksquare$

Examples:

(1) Bern(p) has c.f. $pe^{it} + q$ where $q = 1-p$.

Bin(p) $\stackrel{d}{=} X_1 + \dots + X_n$ where $X_k \sim \text{iid Bern}(p)$, hence by the proposition, Bin(p) has c.f. $(pe^{it} + q)^n$.

(2) Pois(λ) has c.f. $\sum_{n=0}^{\infty} e^{itk} \frac{e^{-\lambda} \lambda^k}{k!}$
 $= e^{-\lambda} e^{\lambda e^{it}}$
 $= e^{-\lambda + \lambda e^{it}}$

(3) Exp(λ) has pdf $\lambda e^{-\lambda x}$ on $(0, \infty)$ and hence its c.f. is $\int_0^{\infty} \lambda e^{-\lambda x} e^{itx} dx$
 $= \lambda \int_0^{\infty} e^{-(\lambda - it)x} dx$
 $= \frac{\lambda}{\lambda - it}$ (why?)

$\Gamma(v, \lambda) \stackrel{d}{=} X_1 + \dots + X_v$ where $X_k \sim \text{iid Exp}(\lambda)$

when v is an integer

and hence, in that case, $\Gamma(v, \lambda)$ has c.f. equal to $\left(\frac{\lambda}{\lambda - it}\right)^v$.

This is also valid for non-integer v ($v > 0$ is always assumed)

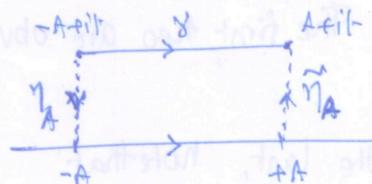
(4) $N(\mu, \sigma^2)$. If $X \sim N(0,1)$ then $\mu + \sigma X \sim N(\mu, \sigma^2)$
Hence $\psi_{N(\mu, \sigma^2)}(t) = e^{i\mu t} \psi_{N(0,1)}(\sigma t) \rightarrow$ Enough to find $\psi_{N(0,1)}$.

For $N(0,1)$,

$$\psi(t) = \int_{\mathbb{R}} e^{itx} \frac{e^{-\frac{1}{2}x^2}}{\sqrt{2\pi}} dx = \int_{\mathbb{R}} e^{-\frac{1}{2}(x - it)^2} \frac{dx}{\sqrt{2\pi}}$$

$$= \lim_{A \rightarrow \infty} \int_{-A}^{+A} e^{-\frac{1}{2}z^2} dz$$

Contour integral along the curve $\gamma_A(u) = A + it + u, u \in [-A, A]$.



By Cauchy's theorem

$$\int_{\gamma_A} e^{-z^2/2} dz = \int_{-A}^A e^{-x^2/2} dx + \int_{\gamma_A} + \int_{\tilde{\gamma}_A}$$

Exer: Show $\int_{\gamma_A}, \int_{\tilde{\gamma}_A} \rightarrow 0$ as $A \rightarrow \infty$

Thus conclude that

$$\psi(t) = e^{-t^2/2}$$

and also, $\psi(t) = e^{i\mu t - \frac{1}{2}\sigma^2 t^2}$
 $N(\mu, \sigma^2)$

Remark: $\psi(t) = e^{-\frac{1}{2}\sigma^2 t^2}$
 $N(0, \sigma^2)$
 $N(0, \sigma^2)$ density = $\frac{e^{-\frac{x^2}{2\sigma^2}}}{\sqrt{2\pi}}$

If $\sigma \downarrow$, the density localizes,
 but ψ becomes more spread out.
 This is an instance of the 'uncertainty principle'.

B) Proposition: If $\hat{\mu} = \hat{\nu}$ then $\mu = \nu$.

We ~~need to~~ ^{shall} use

Parseval's identity: If $\mu, \nu \in \mathcal{P}(\mathbb{R})$, then

$$\textcircled{1} E_{\mu} \psi_{\nu} = E_{\nu} \psi_{\mu}$$

i.e., $\int \psi_{\nu}(t) d\mu(t) = \int \psi_{\mu}(t) d\nu(t)$

$$\textcircled{2} \int e^{-i\alpha t} \psi_{\nu}(t) d\mu(t) = \int \psi_{\mu}(t - \alpha) d\nu(t).$$

Proof: Let $f(s, t) = e^{ist} : \mathbb{R}^2 \rightarrow \mathbb{C}$

Clearly f is integrable w.r.t the probability measure $\mu \otimes \nu$.

Hence, by Fubini's
 $\int_{\mathbb{R}^2} f(s, t) d(\mu \otimes \nu)(s, t)$

$$\begin{aligned} & \int_{\mathbb{R}} \int_{\mathbb{R}} e^{ist} d\mu(s) d\nu(t) && \int_{\mathbb{R}} \int_{\mathbb{R}} e^{ist} d\nu(t) d\mu(s) \\ & = \int_{\mathbb{R}} \psi_{\mu}(t) d\nu(t) && = \int_{\mathbb{R}} \psi_{\nu}(s) d\mu(s). \quad \blacksquare \end{aligned}$$

(b) - exercise.

Proof of the proposition: Let $\theta_{\sigma} = N(0, \sigma^2)$.

By the fact that

$$\begin{aligned} \psi_{\theta_{\sigma}}(t) &= e^{-\sigma^2 t^2/2} = \frac{\sqrt{2\pi}}{\sigma} \frac{1}{\sqrt{2\pi} \frac{1}{\sigma}} e^{-t^2/2(1/\sigma^2)} \\ &= \frac{\sqrt{2\pi}}{\sigma} \cdot \theta_{\frac{1}{\sigma}}(t) \end{aligned}$$

where we use $\theta_{1/\sigma}$ to denote the density of $N(0, \frac{1}{\sigma^2}) = N(0, \frac{1}{\sigma^2})$.

Apply Parseval's to μ and θ_{σ} to get

$$\begin{aligned} \int \hat{\mu}(t) e^{-i\alpha t} d\theta_{\sigma}(t) &= \int \hat{\theta}_{\sigma}(t - \alpha) d\mu(t) \\ &= \frac{\sqrt{2\pi}}{\sigma} \int \theta_{\frac{1}{\sigma}}(t - \alpha) d\mu(t) \\ &= \frac{\sqrt{2\pi}}{\sigma} \underbrace{(\theta_{\frac{1}{\sigma}} * \mu)(\alpha)} \end{aligned}$$

density of the convolution
 $\theta_{\frac{1}{\sigma}} * \mu$.

Now, as $\sigma \rightarrow 0$, $\frac{1}{\sigma} \rightarrow \infty$ and hence

$$\theta_{\frac{1}{\sigma}} \xrightarrow{d} \delta_0$$

Conclude that $\theta_{\frac{1}{\sigma}} * \mu \xrightarrow{d} \delta_0 * \mu = \mu$
 (Why?)

Hence for any $x = \text{cty. point of } F_{\mu}$,

$$F_{\mu}(x) = \lim_{\sigma \rightarrow 0} \int_{-\infty}^x \theta_{\frac{1}{\sigma}} * \mu(\alpha) d\alpha$$

$$= \lim_{\sigma \rightarrow 0} \int_{-\infty}^x \frac{\sigma}{\sqrt{2\pi}} \int_{\mathbb{R}} \hat{\mu}(t) e^{-i\alpha t} d\theta_{\sigma}(t) d\alpha$$

= a function of $\hat{\mu}$ alone!

Thus we can recover $F_{\mu}(x)$ from $\hat{\mu}$
 (for a.e. x and hence for all x , by right continuity of F_{μ}).

Thus $\hat{\mu}$ determines μ . \blacksquare